## $\beta$ -delayed proton decays and spin assignments for <sup>133</sup>Sm and <sup>149</sup>Yb

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Received: 29 May 2006 / Revised: 24 August 2006 / Published online: 11 September 2006 – © Società Italiana di Fisica / Springer-Verlag 2006 Communicated by R. Krücken

**Abstract.** The proton-rich isotope <sup>133</sup>Sm was produced via the fusion evaporation reaction <sup>40</sup>Ca + <sup>96</sup>Ru. Its  $\beta$ -delayed proton decay was studied by p- $\gamma$  coincidence in combination with a He-jet tape transport system, and half-lives, proton energy spectra,  $\gamma$ -transitions following the proton emission, as well as  $\beta$ -delayed proton branching ratios to the low-lying states in the grand-daughter nucleus were determined. Comparing the observed  $\beta$ -delayed proton branching ratios with statistical model calculations, the best agreement is found assuming that only one level with the spin of 3/2 in <sup>133</sup>Sm decays or two levels with the spins of 1/2 and 5/2 decay with similar half-lives. The configuration-constrained nuclear potential energy surfaces of <sup>133</sup>Sm were calculated using the Woods-Saxon-Strutinsky method, which suggests a 1/2<sup>-</sup> ground state and a 5/2<sup>+</sup> isomer with an excitation energy of 120 keV. Therefore, the simple (EC+ $\beta^+$ ) decay scheme of <sup>133</sup>Sm in Eur. Phys. J. A **11**, 277 (2001) has been revised. In addition, our previous experimental data on the  $\beta$ -delayed proton decay of <sup>149</sup>Yb reported in Eur. Phys. J. A **12**, 1 (2001) was also analyzed using the same method. The spin-parity of <sup>149</sup>Yb is suggested to be  $1/2^-$ .

**PACS.** 23.40.Hc Relation with nuclear matrix elements and nuclear structure – 21.10.Hw Spin, parity, and isobaric spin – 24.10.Pa Thermal and statistical models – 27.60 + j 90  $\leq A \leq 149$ 

The  $\beta$ -delayed proton ( $\beta$ p) decay of <sup>133</sup>Sm was first studied by Bogdanov et al. [1] at Dubna in 1977. Its halflife was determined to be  $3.2 \pm 0.4$  s. The spin and parity of  $^{133}$ Sm was suggested to be  $5/2^+$  only based on fitting the measured energy spectrum of  $\beta$ -delayed protons with statistical-model calculations. Experimental data on  $\beta$ p decay of <sup>133</sup>Sm measured by Wilmarth *et al.* [2] using an ISOL facility at LBL was reported in 1985. The half-life of  $^{133}$ Sm was determined to be  $2.8 \pm 0.2$  s. A 213 keV  $\gamma$ -ray following the  $\beta$ p decay of <sup>133</sup>Sm was observed, which corresponds to the transition between the lowest-energy  $2^+$  state to  $0^+$  ground state in the proton daughter nucleus <sup>132</sup>Nd. It was pointed out in [2] that according to a statistical-model calculation the branching ratio to 4<sup>+</sup> state in <sup>132</sup>Nd should be 12%, if the spin-parity of <sup>133</sup>Sm were  $5/2^+$ . However, a 389 keV  $\gamma$ -ray, *i.e.* the  $4^+ \rightarrow 2^+$  transition in <sup>132</sup>Nd with an appropriate intensity was not seen. Therefore, the spin-parity assignment of  $5/2^+$  to <sup>133</sup>Sm is less certain. On the other hand, only the two  $\gamma$ -rays of 369.6 and 156.8 keV with a halflife of  $3.7 \pm 0.7$  s were assigned to the (EC+ $\beta^+$ ) decay by Breitenbach et al. [3] in 1993. Later a simple  $(EC+\beta^+)$ 

decay scheme was proposed by our group [4] in 2001. According to the level scheme of  $^{133}$ Pm given by Regan *et* al. [5], the 214.5 keV + X, 84.5 keV + X and 0.0 keV + X energy levels in our simple decay scheme (see fig. 1 in [4]) were assigned as  $9/2^+$ ,  $7/2^+$  and  $5/2^+$ , respectively. However, the level scheme of <sup>133</sup>Pm was improved and revised by Galindo-Uribarri et al. [6], and the spins and parities of the above three energy levels were assigned as  $7/2^+$ ,  $5/2^+$  and  $3/2^+$  instead in ref. [6]. Of course, a similar revision should be made in our simple decay scheme. It should be noted that the  $\beta$ -delayed  $\gamma$  lines in our simple decay scheme could be separated into two isolated groups. The first group, including the 84.5 keV  $\gamma$  line, which corresponds to the  $5/2^+ \rightarrow 3/2^+$  transition in the daughter nucleus  $^{133}$ Pm, could be from a  $5/2^+$  state of  $^{133}$ Sm with a half-life of  $2.8 \pm 0.5$  s. The half-life of another group which includes the 369.6 and 156.8 keV  $\gamma$ -rays was determined to be  $3.4 \pm 0.5$  s, and seems different from the first one.

In the present work, a new study on  $\beta$ p decay of <sup>133</sup>Sm is reported, and the spin-parity of <sup>133</sup>Sm is proposed. Comparing the  $\beta$ p decay and  $\beta$ -delayed  $\gamma$  decay, the simple (EC+ $\beta^+$ ) decay scheme of <sup>133</sup>Sm in [4] is revised. In addition, our previous experimental data on  $\beta$ p decay of

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Fig. 1. The measured  $\gamma$ - and X-ray spectrum in coincidence with delayed protons of 2.5 to 6.0 MeV in the reaction  ${}^{40}\text{Ca} + {}^{96}\text{Ru}$ . The intense peaks are labeled by their energies in keV and their  $\beta$ p precursors.

 $^{149}{\rm Yb}$  [7] was analyzed with statistical-model calculations, and a spin-parity of  $^{149}{\rm Yb}$  is proposed.

The experiment described here was carried out at the Sector-Focusing Cyclotron in the Institute of Modern Physics, Lanzhou, PRC. A schematic view of the experimental set-up is shown in fig. 1 of ref. [8]. A 232 MeV  ${}^{40}\text{Ca}^{12+}$  beam entered a target chamber filled with 1 bar helium through a  $1.89 \,\mathrm{mg/cm^2}$  thick Havar window. After traversing a layer of helium gas and an aluminium degrader, the beam finally hit on a  $^{96}$ Ru target (85% enriched) with a thickness of about  $1.4 \,\mathrm{mg/cm^2}$ . The beam energy in the middle of the target was 180 MeV, and the beam intensity was about 40 pnA. The <sup>133</sup>Sm was produced via the 2pn evaporation channel. We used a heliumjet in combination with a tape transport system to periodically move the radioactivity into a shielded counting room, using  $PbCl_2$  at 430 °C as aerosol. The length of the capillary is about 6 m. The collection time, tape moving time, waiting time, and accumulation time was adjusted to 1.60, 0.16, 0.16, and 1.44s, respectively. In order to study the  $\beta p$  decay, proton-gamma coincidence measurements were carried out [8–10]. Two  $570 \,\mathrm{mm^2} \times 350 \,\mu\mathrm{m}$ , totally depleted silicon surface barrier detectors located on opposite sides of the movable tape were used to detect the protons. Behind each silicon detector a coaxial HpGe(GMX) was placed to observe  $\gamma$ - and X-ray. The energy and time spectra of  $\gamma$ - and X-rays as well as protons were recorded in both single and coincidence modes.

The observed  $\gamma$ - and X-ray spectrum gated on  $\beta$ delayed protons of 2.5–6.0 MeV is shown in fig. 1. The intense peaks, except 511 keV and X-rays, are labeled with their energies in keV and their  $\beta$ p precursors. In particular, the most intense 213 keV  $\gamma$  line was assigned to the transition between the lowest-energy  $2_1^+$  state and  $0^+$ ground state in the proton daughter nucleus <sup>132</sup>Nd [11] following the  $\beta$ p decay of <sup>133</sup>Sm, while the 398, 611 and 824 keV  $\gamma$  lines were assigned to the  $4^+ \rightarrow 2_1^+, 2_2^+ \rightarrow 2_1^+$ and  $2_2^+ \rightarrow 0^+$  transitions in <sup>132</sup>Nd [11], respectively.

The energy spectrum of  $\beta$ -delayed protons gated on the 213 keV  $\gamma$  line is shown in fig. 2, while the inset displays



Fig. 2. The energy spectrum of  $\beta$ -delayed protons gated on the 213 keV  $\gamma$  line in <sup>133</sup>Sm. The histogram is the experimental result, while the curves stand for the statistical-model calculations. The inset displays the decay curve of the 213 keV line coincident with  $\beta$ -delayed protons.

**Table 1.** Calculated and experimental relative branching ratios  $(b_{\beta p})$  to different final states in <sup>132</sup>Nd from the  $\beta p$  decay of <sup>133</sup>Sm, assuming different values of the initial spin and parity of <sup>133</sup>Sm. The experimental relative  $b_{\beta p}$  leading to the  $2_1^+$  state was normalized to 50.

Initial spin	Relative $b_{\beta p}$ to the final states (%)			
and parity	G. S.	$213\mathrm{keV}$	$611\mathrm{keV}$	$824\mathrm{keV}$
of $^{133}$ Sm	$(0^+)$	$(2_1^+)$	$(4^{+})$	$(2_2^+)$
$1/2^{-}$	56.6	39.6	0.5	3.3
$1/2^+$	51.3	43.5	1.1	4.0
$3/2^{-}$	37.7	54.0	3.7	4.7
$3/2^+$	38.3	54.1	2.4	5.2
$5/2^{-}$	22.5	60.9	11.1	5.5
$5/2^{+}$	15.8	64.2	13.7	6.3
Exp. Value		$50 \pm 4$	$3.0\pm0.5$	$5\pm1$

the decay curve of the 213 keV  $\gamma$  line when gated on the  $\beta$ -delayed protons. From the decay curve the "average" half-life of <sup>133</sup>Sm was extracted to be  $3.2 \pm 0.7$  s.

The relative branching ratios  $(b_{\beta p})$  to different final states in the proton daughter nucleus <sup>132</sup>Nd observed in the  $\beta$ p decay of <sup>133</sup>Sm for various values of the initial spin and parity of <sup>133</sup>Sm and the proton energy spectra were calculated using a revised statistical model [12,13]. The calculated results are listed in table 1 together with the experimental values, which were determined by means of the relative intensities of the correspondent  $\gamma$  lines in fig. 1. Unfortunately, we could not obtain the experimental  $b_{\beta p}$ to the ground state in <sup>132</sup>Nd because of the restriction of the p- $\gamma$  coincidence method. Comparing the calculated branching ratios with the experimental values, the closest agreement is found assuming that only one level with the spin of 3/2 in <sup>133</sup>Sm decays or two levels with the spins of 1/2 and 5/2 decay with similar half-lives.



Fig. 3. Calculated nuclear potential energy surfaces of  $1/2^-$  ground state (a) and  $5/2^+$  isomer (b) for <sup>133</sup>Sm.



**Fig. 4.** Revised simple  $(EC+\beta^+)$  decay scheme of <sup>133</sup>Sm.

To shed more light on the ground-state properties of  $^{133}$ Sm, the configuration-constrained nuclear potential energy surface (NPES) (fig. 3) was calculated by using the Woods-Saxon-Strutinsky method [14]. If the potential energy of <sup>133</sup>Sm is minimum, its shape has to be prolate, and the last neutron in <sup>133</sup>Sm is not able to occupy the obital of  $\nu 3/2^+$ [402]. Actually, a minimum at deformation parameters  $\beta_2 = 0.327$  and  $\gamma = -0.005^{\circ}$  was found, which corresponds to the configuration  $\nu 1/2^{-541}$  other than  $\nu 3/2^+[402]$ , and the second found out minimum at  $\beta_2 = 0.291$  and  $\gamma = 0.000^\circ$  corresponds to the configuration of  $\nu 5/2^+$  [402]. The excitation energy of the minimum for the configuration of  $\nu 5/2^+$  [402] is only 120 keV higher than that for the configuration of  $\nu 1/2^{-}$ [541]. In addition, the calculated energy spectra of  $\beta$ -delayed protons fit the experimental data well assuming two decaying levels in <sup>133</sup>Sm with spins and parities  $1/2^-$  and  $5/2^+$  (see fig. 2). Therefore, the spins and parities of the two decaying levels in  $^{133}$ Sm were suggested to be  $1/2^-$  (ground state) and  $5/2^+$  (isomeric state). Comparing the calculated  $b_{\beta p}$ and the experimental values, the relative decay intensities of  $1/2^-$  state and  $5/2^+$  state are ~ 75% and ~ 25%, respectively. Finally, the simple  $(EC+\beta^+)$  decay scheme of  $^{1\overline{3}3}$ Sm in [4] should be separated into two components and revised in fig. 4. The first component of the  $(EC+\beta^+)$  decay with the half-life of 2.8 s, including the 84.5 keV  $\gamma$  line,

**Table 2.** Calculated and experimental relative branching ratios  $(b_{\beta p})$  to different final states in <sup>148</sup>Er from the  $\beta$ p decay of <sup>149</sup>Yb, assuming different values of the initial spin and parity of <sup>149</sup>Yb. The experimental relative  $b_{\beta p}$  leading to the 2<sup>+</sup> state was normalized to 50.

Initial spin	Relative $b_{\beta p}$ to the final states (%)			
and parity	G. S.	$647\mathrm{keV}$	$1524\mathrm{keV}$	
of $^{149}$ Yb	$(0^+)$	$(2^+)$	$(4^{+})$	
$1/2^{-}$	83.0	16.8	0.2	
$1/2^+$	77.7	22.1	0.2	
$3/2^{-}$	70.0	29.3	0.7	
$3/2^+$	70.7	28.8	0.5	
$5/2^{-}$	55.3	42.2	2.4	
$5/2^{+}$	49.8	47.2	3.0	
Exp. Value		50	≤ 1	

comes from the 5/2<sup>+</sup> isomeric state in <sup>133</sup>Sm, while the second one with a half-life of 3.4 s, including the 369.6 and 156.8 keV  $\gamma$  lines, comes from the 1/2<sup>-</sup> ground state. By the way, the "average" half-life of the 213 keV  $\gamma$  line gated on  $\beta$ -delayed protons is 3.2 s, somewhere between 2.8 and 3.4 s, because of the contributions of the two components.

The experimental data of  $\beta p$  decay of <sup>149</sup>Yb was reported for the first time by our group [7] in 2001. The relative  $b_{\beta p}$  to different final states in the proton daughter nucleus <sup>148</sup>Er from the  $\beta p$  decay of <sup>149</sup>Yb for various values of the initial spin and parity of <sup>149</sup>Yb were calculated with the revised statistical model. The calculated  $b_{\beta p}$  are listed in table 2 together with the experimental values. Comparing the calculated  $b_{\beta p}$  with the experimental values, the initial spin and parity of <sup>149</sup>Yb could be 1/2 or 3/2. Recently, the configuration-constrained NPES of <sup>149</sup>Yb was calculated using the Woods-Saxon-Strutinsky method. A minimum with  $\beta_2 = 0.166$  and  $\gamma = 60^{\circ}$  was found (fig. 5), which corresponds to the configuration of  $\nu 1/2^{-}$ [521]. In addition, the  $\beta p$  energy spectrum of <sup>149</sup>Yb



Fig. 5. Calculated nuclear potential energy surface for <sup>149</sup>Yb.



Fig. 6. The energy spectrum of  $\beta$ -delayed protons gated on the 647 keV  $\gamma$  line in <sup>149</sup>Yb. The histogram is the experimental result, while the curve stands for the statistical-model calculation with initial spin and parity of  $1/2^{-}$ .

gated on the 647 keV  $\gamma$ -ray was calculated using the revised statistical model, assuming the spin and parity of <sup>149</sup>Yb to be  $1/2^-$ . The calculated spectrum fitted the experimental data very well (see fig. 6). Finally, the ground-state spin-parity of <sup>149</sup>Yb is proposed to be  $1/2^-$ .

This work was supported by the National Natural Science Foundation of China (10375078 and 10475002).

## References

- D.D. Bogdanov, A.V. Demyanov, V.A. Karnaukhov *et al.*, Nucl. Phys. A **275**, 229 (1977).
- P.A. Wilmarth, J.M. Nitschke, J.M. Lemmertz *et al.*, Z. Phys. A **321**, 179 (1985).
- J. Breitenbach, R. Braga, J.L. Wood et al., in 6th International Conference on Nuclei Far From Stability / 9th International Conference on Atomic Masses and Fundamental Constants, 1992, Inst. Phys. Conf. Series, No. 132, 575 (1993).
- S.-W. Xu, Y.-X. Xie, X.-D. Wang *et al.*, Eur. Phys. J. A 11, 277 (2001); 12, 375 (2001) (E).
- 5. P.H. Regan et al., Nucl. Phys. A 533, 476 (1991).
- A. Galindo-Uribarri, D. Ward, H.R. Andrews *et al.*, Phys. Rev. C 54, 1057 (1996).
- S.-W. Xu, Z.-K. Li, Y.-X. Xie *et al.*, Eur. Phys. J. A **12**, 1 (2001).
- S.-W. Xu, Z.-K. Li, Y.-X. Xie *et al.*, Phys. Rev. C 71, 054318 (2005).
- S.-W. Xu, Y.-X. Xie, Z.-K. Li *et al.*, Phys. Rev. C **60**, 061302(R) (1999).
- S.-W. Xu, Y.-X. Xie, Z.-K. Li *et al.*, Z. Phys. A **356**, 227 (1996).
- Yu Khazov, A.A. Rodinov, S. Sakharov *et al.*, Nucl. Data Sheets **104**, 497 (2005).
- P. Hornshoj, K. Wilsky, P.G. Hansen *et al.*, Nucl. Phys. A 187, 609 (1972).
- 13. J.C. Hardy, Phys. Lett. B 109, 242 (1982).
- W. Nazarewicz, J. Dudek, R. Bengtsson *et al.*, Nucl. Phys. A **435**, 397 (1985).